

Geometric and Electronic Structure of the Mn(IV)Fe(III) Cofactor in Class Ic Ribonucleotide Reductase: Correlation to the Class Ia Binuclear Non-Heme Iron Enzyme

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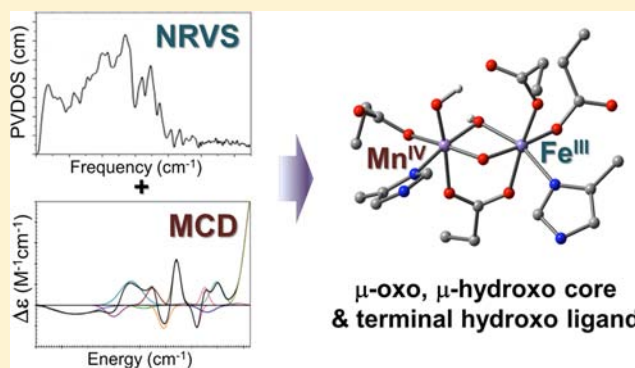
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Supporting Information

ABSTRACT: The class Ic ribonucleotide reductase (RNR) from *Chlamydia trachomatis* (*Ct*) utilizes a Mn/Fe hetero-binuclear cofactor, rather than the Fe/Fe cofactor found in the β (R2) subunit of the class Ia enzymes, to react with O₂. This reaction produces a stable Mn^{IV}Fe^{III} cofactor that initiates a radical, which transfers to the adjacent α (R1) subunit and reacts with the substrate. We have studied the Mn^{IV}Fe^{III} cofactor using nuclear resonance vibrational spectroscopy (NRVS) and absorption (Abs)/circular dichroism (CD)/magnetic CD (MCD)/variable temperature, variable field (VTVF) MCD spectroscopies to obtain detailed insight into its geometric/electronic structure and to correlate structure with reactivity; NRVS focuses on the Fe^{III}, whereas MCD reflects the spin-allowed transitions mostly on the Mn^{IV}. We have evaluated 18 systematically varied structures. Comparison of the simulated NRVS spectra to the experimental data shows that the cofactor has one carboxylate bridge, with Mn^{IV} at the site proximal to Phe₁₂₇. Abs/CD/MCD/VTVF MCD data exhibit 12 transitions that are assigned as d–d and oxo and OH[−] to metal charge-transfer (CT) transitions. Assignments are based on MCD/Abs intensity ratios, transition energies, polarizations, and derivative-shaped pseudo-A term CT transitions. Correlating these results with TD-DFT calculations defines the Mn^{IV}Fe^{III} cofactor as having a μ -oxo, μ -hydroxo core and a terminal hydroxo ligand on the Mn^{IV}. From DFT calculations, the Mn^{IV} at site 1 is necessary to tune the redox potential to a value similar to that of the tyrosine radical in class Ia RNR, and the OH[−] terminal ligand on this Mn^{IV} provides a high proton affinity that could gate radical translocation to the α (R1) subunit.



1. INTRODUCTION

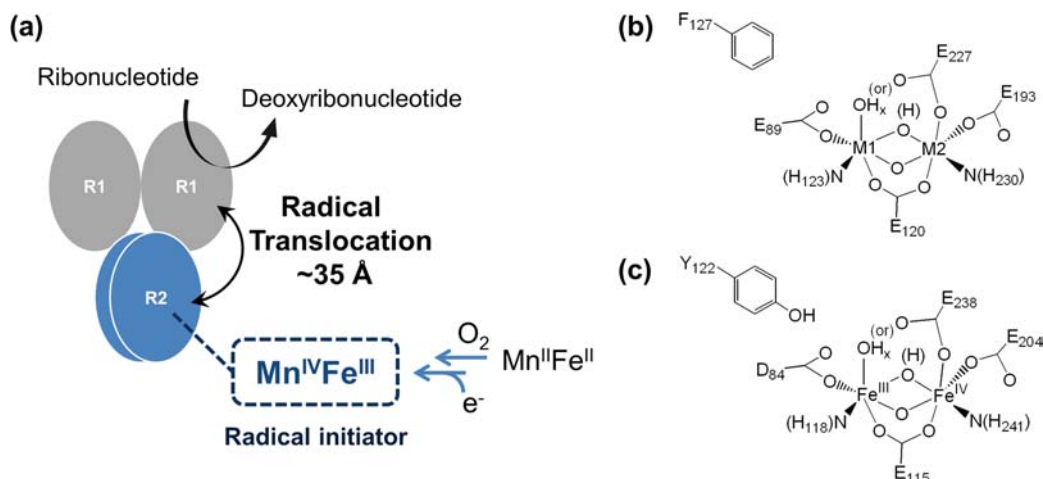
Ribonucleotide reductase (RNR) plays a key role in DNA synthesis by reducing ribonucleotides to their corresponding 2'-deoxyribonucleotides.^{1–3} Among the three different classes of RNRs, the class I enzymes have 2 subunits, α (R1) and β (R2). The α (R1) subunit directly reduces the substrates (ribonucleotides) using radical chemistry (thiyl radical with disulfide bond formation), while the β (R2) subunit initiates an electron hole by the reaction of a binuclear nonheme metal active site with O₂.

In the class Ia RNRs, the electron hole is generated by the high-valent intermediate, X, which abstracts a hydrogen atom (H-atom) from a nearby tyrosine.² This tyrosine radical then transfers its electron hole over ~ 35 Å by a series of hole-

hopping and/or proton-coupled electron-transfer (PCET) steps,^{4,5} ultimately leading to the initiation of substrate reduction by abstraction of the 3'-H. A number of structural studies on the intermediate X have been performed using Mössbauer, electron paramagnetic resonance (EPR), electron–nuclear double resonance (ENDOR), extended X-ray absorption fine structure (EXAFS), and magnetic circular dichroism (MCD) spectroscopies.^{6–18} Based on these studies, X was shown to have an Fe^{III}Fe^{IV} center that has strong antiferromagnetic coupling between two high-spin irons, giving an $S = 1/2$ ground state reflecting a bridged structure with at

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Scheme 1^a

^a(a) Function of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor in catalysis by a class Ic ribonucleotide reductase: substrate reduction in α (R1) following radical initiation by $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor in β (R2) through a radical-translocation process; (b) structure of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor site in class Ic ribonucleotide reductase ($x = 1$ or 2 , M1 = metal site 1, and M2 = metal site 2); and (c) structure of the $\text{Fe}^{\text{III}}\text{Fe}^{\text{IV}}$ intermediate X in class Ia ribonucleotide reductase ($x = 1$ or 2).

least one oxo-bridge.⁹ The Fe^{IV} center is at the distal position to the tyrosine residue (Y_{122} , Scheme 1c),^{16,18} which is oxidized by X to produce a stable tyrosine radical. These studies have provided fundamental insight into the O_2 reaction mechanism of the biferrrous site and how the tyrosine radical ($\text{Y}\bullet$) is generated. However, there is still an issue concerning whether one oxo bridge is protonated.^{13b,16,17,19}

In the class Ic RNRs, the position occupied by the tyrosine residue (Y_{122}) in the class Ia (and Ib) orthologues is instead occupied by a phenylalanine (F_{127}), which does not generate a radical.²⁰ A number of studies have determined that in the class Ic RNRs, the metal center with the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ redox state is directly responsible for the radical initiation.^{21–24} This oxidized cofactor is generated by the reaction of the $\text{Mn}^{\text{II}}\text{Fe}^{\text{II}}$ site with O_2 and proceeds through a $\text{Mn}^{\text{IV}}\text{Fe}^{\text{IV}}$ intermediate that undergoes one-electron reduction (via a residue at the surface of the protein)²⁵ to form the catalytically active $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ site (Scheme 1a,b).²⁶ Mössbauer, extended X-ray absorption fine structure (EXAFS), and anomalous diffraction crystallography have been employed to study this cofactor site.^{22,27–29} This $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ center has an $S = 1$ ground state from strong antiferromagnetic coupling between the Mn and Fe centers.²² A structure with oxo, hydroxo bridges was favored based on the atomic distances,²⁷ and the Mn was observed in the metal site 1 (close to F_{127} , Scheme 1b).^{28–30} However, a range of questions still remains in terms of its geometric and electronic structure and how its electron–hole initiation and transport mechanistically relate to the analogous processes in the tyrosine-radical-dependent class Ia orthologs.

The goal of this study is to further define the geometric and electronic structure of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor site in class Ic RNRs, which is important as a functional analog of the radical initiator $\text{Y}\bullet$ as well as a structural analog of the high-valent intermediate X in the class Ia enzymes. In this study, a suite of spectroscopic approaches, including absorption (Abs), circular dichroism (CD), magnetic CD (MCD), variable temperature, variable field (VT-VH) MCD, and nuclear resonance vibrational spectroscopies (NRVS), has been used to define the geometric and electronic structure of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor in *Ct* RNR. NRVS is a relatively new, third-generation synchrotron-based

technique that measures the inelastic vibrational side bands of the ^{57}Fe Mössbauer nuclear transition at 14.4 keV that involve Fe motion^{31,32} and is therefore a direct probe of the Fe site. MCD focuses on ligand-field (LF) and charge-transfer (CT) transitions.³³ High MCD to Abs intensity is associated with spin-allowed LF transitions of the Mn^{IV} (as high-spin Fe^{III} LF transitions are all spin-forbidden), which provide a direct probe of the Mn site, whereas the CT transitions probe the ligands bound to both the Mn^{IV} and the Fe^{III} . Coupled with DFT calculations, this study defines the geometric and electronic structure of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ active site in *Ct* RNR and its correlation to the binuclear nonheme iron site in the class Ia enzymes.

2. EXPERIMENTAL SECTION

2.1. Sample Preparation. Samples of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ complex of *Ct* RNR were prepared as previously described.²² The MCD samples were made by mixing the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ complex in buffer (100 mM HEPES, H_2O or D_2O) with glycerol (or glycerol-*d*₃, 50% v/v for final concentration) as a glassing agent for low-temperature spectral measurements. Samples of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ complex for NRVS were prepared with ^{57}Fe , and the species present in the samples were quantified by Mössbauer spectroscopy (Figures S1 and S2). The samples were injected into custom-made NRVS cells and frozen in liquid N_2 .

2.2. Spectroscopic Methods. The room-temperature Abs spectrum was measured on an Agilent 8453 UV–vis spectrophotometer. The low-temperature Abs spectrum was obtained by using a Cary 500 spectrophotometer with a Janis STVP-100 continuous-flow liquid helium (He) cryostat equipped with a Lakeshore model 332S temperature controller.

UV–vis (300–900 nm region) CD and MCD data were collected on a Jasco J810 spectrometer with an extended S-20 photomultiplier tube, and near-infrared (NIR, 600–2000 nm region) CD and MCD data were collected on a Jasco J730 spectrometer with a liquid N_2 -cooled InSb detector. The MCD spectra were obtained at temperatures between 1.8 and 50 K. The sample temperature was measured by using a calibrated Cernox resistor (Lakeshore Cryogenics) inserted into the MCD cell. Magnetic fields from 0 to 7 T were applied by using an Oxford magnet (either SM 4-7T or SM 4000-7T). For each field and temperature, (+) and (–) signals were averaged for baseline correction and to improve the signal-to-noise (S/N) ratios. VTVH

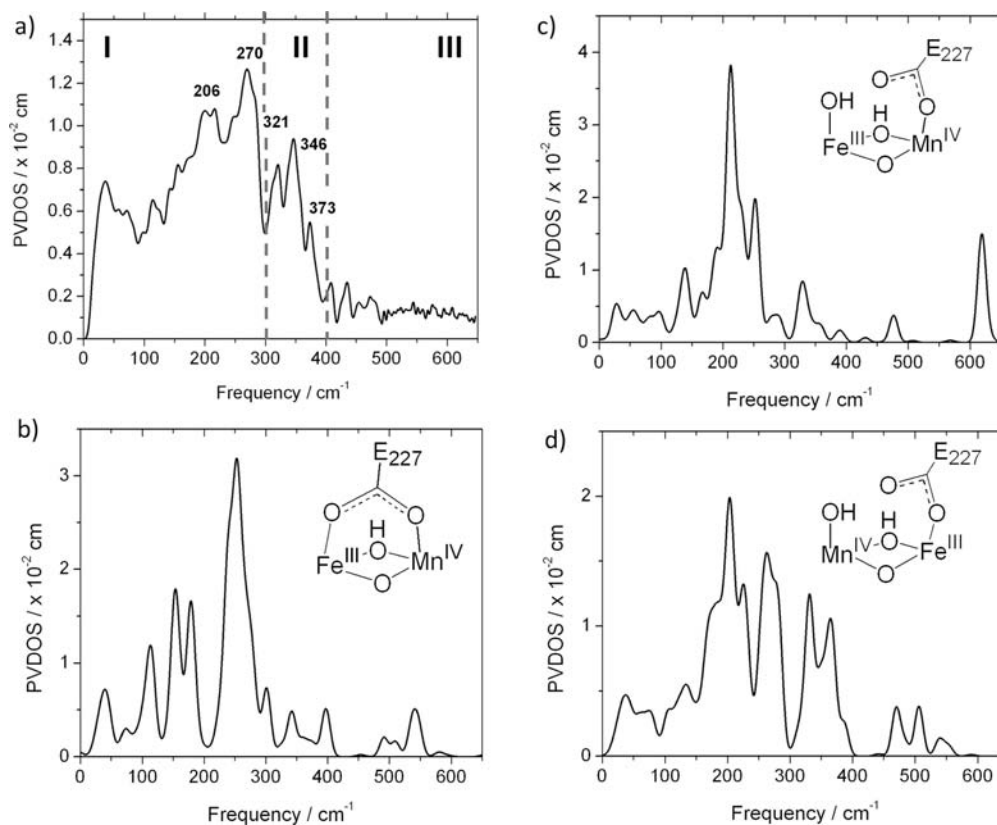


Figure 1. (a) NRVS spectra of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor and representative simulations for structures having (b) two carboxylate bridges (bridging E_{227}); (c) one carboxylate bridge (terminal E_{227}) with Fe^{III} at site 1; and (d) one carboxylate bridge (terminal E_{227}) with Mn^{IV} at site 1.

MCD data were normalized to the observed intensity at lowest temperature and highest field (1.8 K, 7 T) over all isotherms for a given wavelength. The VTVH MCD data were fit to extract values of the spin Hamiltonian, with the ground-state parameters given by the spin projection of the single metal site into the dimer wave functions (see ref 33).

NRVS spectra were collected at beamline BL09XU at SPring-8 in Hyogo, Japan and at beamline 3-ID-D at the Advanced Photon Source, Argonne, U.S.A. on multiple visits. Spectra were collected between -25 and 80 meV in 0.25 meV steps. During data collection, the samples were maintained at 20 – 30 K by using a liquid He cryostat. Photoreduction effects were evaluated by comparing successive spectra during data acquisition. Scans were averaged to obtain better S/N ratios. Averaged data were converted to partial-vibrational density-of-states (PVDOS) spectra using the PHOENIX program.³⁴

2.3. Computational Methods. The starting geometry was obtained from the crystal structure of Fe-loaded *Ct* RNR R2 (1SYR RCSB structure).²⁰ The protein active site residues were truncated by replacing the β -carbons with methyl groups. The positions of these β -carbons were fixed to impose the protein backbone constraints for the geometry optimization but not the frequency calculations. This allowed for the large number of calculations required and gave relative energies very similar to those obtained in calculations in which the α -carbons were constrained.^{35,36} Calculations were performed with the Gaussian 09 software package³⁷ unless specified otherwise.

For NRVS spectral simulations, geometry optimizations and frequency calculations were performed using the spin-unrestricted BP86 functional with the LANL2DZ basis set. Previous studies on inorganic complexes have shown that this combination reproduced NRVS spectra for known structures.^{38–41} The vibrational modes were converted to NRVS PVDOS spectra using the *gennrvs* program.³⁸ These spectra were further validated with parallel-simulated spectra using heavy masses at the β -carbons and their hydrogens to limit the effect of protein backbone motion.

For LF and CT transitions, geometry optimizations and TD-DFT calculations were performed by using the spin-unrestricted B3LYP functional with the 6-311+g* basis set on the metals and 6-31g* on the rest of the atoms. This functional and basis set combination was calibrated by comparison to data on inorganic complexes, as described in Section 3.2.4.1. Zero-field splitting (ZFS) parameter D and E were calculated for these optimized structures by using the Orca 2.8.0 program^{42,43} with the B3LYP functional. The CP(PPP) basis set was used for the metals, and the TZVP basis set was used for the rest of the atoms.

Total energies of the different candidate structures were calculated with the spin-unrestricted B3LYP* functional (15% Hartree–Fock (HF) exchange rather than the conventional 20%) and the 6-311+g* basis set for the geometry-optimized structures (spin-unrestricted BP86 functional with LANL2DZ basis set) for comparison to previous calculations.⁴⁴ The effects of the protein environment were included by using the polarizable continuum model (PCM) with a dielectric constant (ϵ) of 4.0. Calculations of exchange coupling constants were performed by using the Orca 2.8.0 program, with the spin-unrestricted BP functional and the def2-SVP basis set. Mössbauer parameters were calculated by using the Orca 2.8.0 program, with the spin-unrestricted BP functional. The CP(PPP) basis set was used for the metals, and DZP basis set was used for the rest of the atoms. The parameters obtained were corrected empirically on the basis of a correlation between theory and experiment.⁴⁵

3. RESULTS AND ANALYSES

3.1. NRVS. 3.1.1. Spectrum. The partial vibrational density of states (PVDOS) spectrum of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor of *Ct* RNR- β (R2) is shown in Figure 1a. This PVDOS intensity reflects the amount of Fe motion in a molecular vibrational mode at a given energy. On the basis of this selectivity and the correlation with structural variations described below, the experimental NRVS spectrum is divided into three regions:

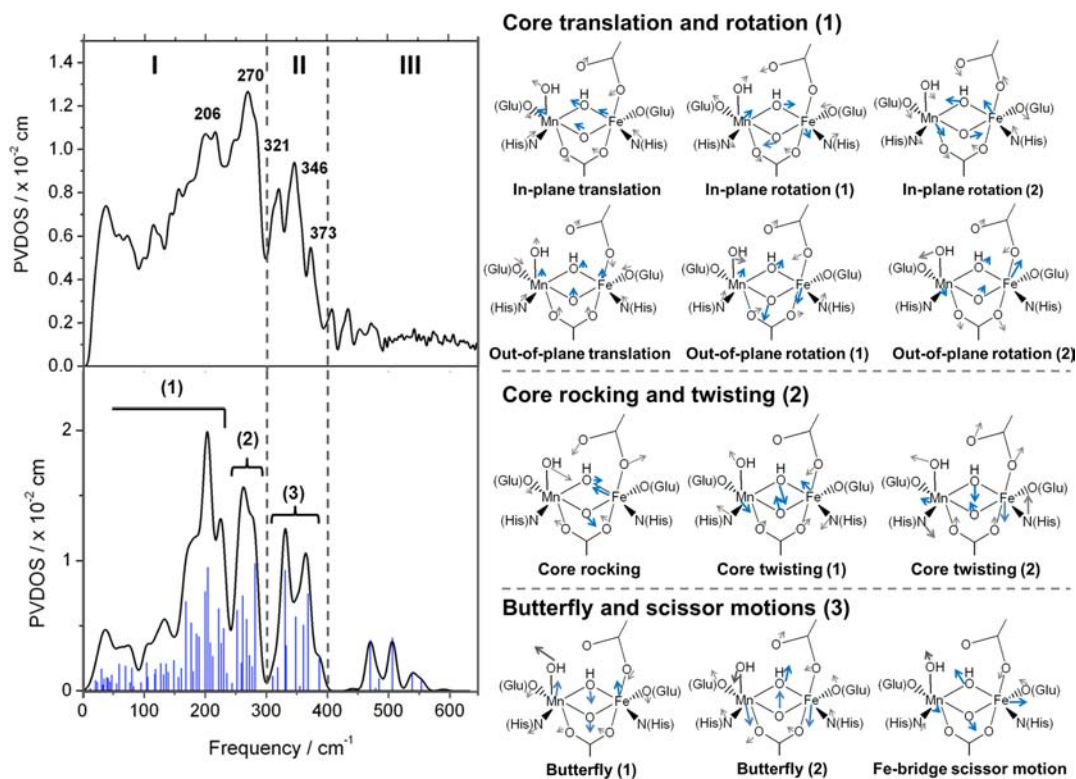


Figure 2. (Left) Spectra obtained from NRVS (experimental, top; simulated, bottom) and (right) related core motions contributing to the features in the simulated spectrum.

intense features in the low-energy ($200\text{--}300\text{ cm}^{-1}$) region (region I), moderately intense features in the $300\text{--}400\text{ cm}^{-1}$ region (region II), and absence of features in the high-energy region above 400 cm^{-1} (region III). To obtain insight into the geometric structure of the complex, this experimental spectrum was compared to a series of NRVS spectra calculated by DFT for different possible structures (Section 3.1.2)

3.1.2. DFT Calculations and Spectral Simulations. In past studies on structurally defined model complexes, the NRVS spectra calculated from DFT-optimized structures correlated well with the experimental NRVS spectra, indicating that NRVS is a powerful technique for defining geometric structure.^{38–41} To analyze the NRVS spectra, DFT calculations were performed for 18 different possible structural models of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor in *Ct* RNR- β (R2). It was assumed that this cofactor has some structural similarity to the $\text{Fe}^{\text{III}}\text{Fe}^{\text{IV}}$ intermediate, X, in *E. coli* RNR- β (R2), and four different structural variations were therefore evaluated: (1) one vs two carboxylate bridges (i.e., terminal bound E_{227} on metal site 2 vs $\mu\text{-}1,3\text{ E}_{227}$ residue), (2) Mn or Fe at metal site 1, which is the metal site proximal to F_{127} that replaces Y_{122} of the class Ia RNR- β (R2) protein from *E. coli* (Scheme 1b), (3) bis($\mu\text{-oxo}$) vs $\mu\text{-oxo}$, $\mu\text{-hydroxo}$ core, and (4) for the structures with a single carboxylate bridged structure (i.e., terminal E_{227} residue at metal site 2), either a terminal water or a hydroxo ligand at metal site 1.

The simulated NRVS spectra for 18 different structures are shown in Figure S3, which are represented by the NRVS spectra shown in Figure 1b–d. Comparison of simulated spectra to the experimental spectrum in Figure 1a provides systematic structural insight. For the structures with two carboxylate bridges, the calculated NRVS spectra show very little NRVS intensity in region II ($300\text{--}400\text{ cm}^{-1}$) as

represented in Figure 1b, regardless of which metal is at site 1 of the core bridging structure (first column of Figure S3). Thus, the presence of the features in region II in the experimental spectrum in Figure 1a requires that the complex possess at most one, and not two, bridging carboxylate ligands. In the cases of single carboxylate bridged structures, the calculated NRVS spectra are very sensitive to the metal (Mn or Fe) at site 1, as shown in Figure 1c,d. For the single-carboxylate bridged structures with Fe at site 1, intense features are calculated to be present in region III ($400\text{--}600\text{ cm}^{-1}$), as shown in Figure 1c, which are absent in the experimental data. While these features in region III are less intense for the structures with a terminal water ligand on an Fe at site 1, this leads to features in region I at energies below 240 cm^{-1} , which do not correlate with the experimental data (Figure S3, second column). Thus the NRVS intensities in region III and in region I below 240 cm^{-1} together, as shown in Figure 1a vs c, rule out all single carboxylate bridged structures with Fe at site 1. For the single-carboxylate bridged structure with Mn at site 1, as shown in Figure 1d, the calculated spectrum agrees well with the experimental NRVS spectrum, having intense features in region I and moderately intense features in region II. For these structures, the high-energy features are calculated to be weak (totaling $<10\%$ of the NRVS intensity distributed over region III) and do not contribute to the NRVS spectrum above the noise level. Therefore, comparisons of the calculated to the experimental spectrum indicate that the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ complex has a single carboxylate bridge (a terminal carboxylate E_{227} on Fe at site 2) and Mn at site 1. These structures could have either a bis($\mu\text{-oxo}$) or $\mu\text{-oxo}$, $\mu\text{-hydroxo}$ core, and either a terminal hydroxo or water on the Mn.

While the calculated NRVS spectra are not very diagnostic of the core bridging structure (bis($\mu\text{-oxo}$) vs $\mu\text{-oxo}$, $\mu\text{-hydroxo}$) or

the nature of the terminal ligand (OH^- vs H_2O) on Mn at site 1, further analyses of these models limit the structural possibilities. The bis(μ -oxo) core structure with a terminal OH^- ligand on the Mn at site 1 significantly lowers the intensity in region II ($300\text{--}400\text{ cm}^{-1}$) as compared to the experimental NRVS spectrum, due to the poor linkage between the Mn and Fe sites (the short Mn- μ -oxo and Mn-terminal OH bonds prevent mixing of Fe motion into the Mn motion) and therefore can be ruled out (Figure S3, right column top spectrum). In addition, among the structures with a μ -oxo, μ -hydroxo core, structures with the hydroxo bridge on the side of the Mn/Fe core with the two His ligands (Figure S3, right column fifth and sixth spectra from the top) were at higher energy (by 4–9 kcal/mol) compared with structures with the hydroxo bridge on the side of the Mn/Fe core with 2 carboxylate ligands (Figure S3, right column third and fourth spectra from the top), consistent with previous computational results.⁴⁴ Therefore, of the six possible structures with a single carboxylate bridge and Mn at site 1, three structures remain viable: a bis(μ -oxo) core with a terminal H_2O ligand; a μ -oxo, μ -hydroxo core with a terminal H_2O ligand; and a μ -oxo, μ -hydroxo core with terminal OH^- ligand.

3.1.3. Spectral Assignments. The previous section narrows the possible structural models for the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor to 3 out of the 18 originally considered (Figure S3, right column, second–fourth spectra from the top). These three remaining structures give similar normal modes with Fe displacements at similar energies. Here, we assign and describe the modes contributing to each region by consideration of the structure with μ -oxo, μ -hydroxo core, and terminal OH on the Mn (Figure 2), which is supported by the MCD results in the next section.

In the low-energy region up to $\sim 300\text{ cm}^{-1}$ (region I), the normal modes involve two types of motions: core translations and rotations (the Mn-(μ -O/ μ -OH)-Fe core moving in the protein pocket, Figure 2, right (1)) and core rocking (Fe moving along the μ -O bridge and bridging atoms moving in the opposite direction) and twisting motions (pairs of metal- μ -O(H) moving in opposite directions, Figure 2, right (2)). For the $300\text{--}400\text{ cm}^{-1}$ region (region II), the features in the NRVS spectrum are mainly attributable to the butterfly motion (metals moving out of the plane with bridging O (and OH) moving in the opposite direction, Figure 2, right (3)), with a contribution from the Fe scissor motion (Fe moving along the metal–metal vector with the bridges moving out of the plane, in the perpendicular direction, Figure 2, right (3)). High-energy features in region III ($>400\text{ cm}^{-1}$) would arise primarily from metal–ligand (OH^- or H_2O) stretching motions for the metal in site 1 having the terminal OH^- or H_2O ligand. These have very low intensity when Mn is at site 1 with a single carboxylate bridge, because the Mn-terminal OH_x stretch involves little Fe motion. This consideration explains the absence of features in region III in the experimental data in Figure 2a.

3.2. Abs/CD/MCD Spectroscopy. **3.2.1. Spectra.** Abs, low-temperature CD and MCD spectra for the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ complex of Ct RNR- β (R2) are shown in Figure 3. The sample contained $\sim 13\text{--}15\%$ $\text{Fe}^{\text{III}}\text{Fe}^{\text{III}}$ and $\sim 7\text{--}10\%$ $\text{Mn}^{\text{III}}\text{Fe}^{\text{III}}$ component, as determined by analysis of the Mössbauer spectrum of a sample prepared in parallel with the NRVS sample (Figures S1 and S2 and Table S1). The contributions of these components to the spectra are negligible in comparison with those from the predominant $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ species (spectra for the pure $\text{Fe}^{\text{III}}\text{Fe}^{\text{III}}$ form are shown in Figure S6). A

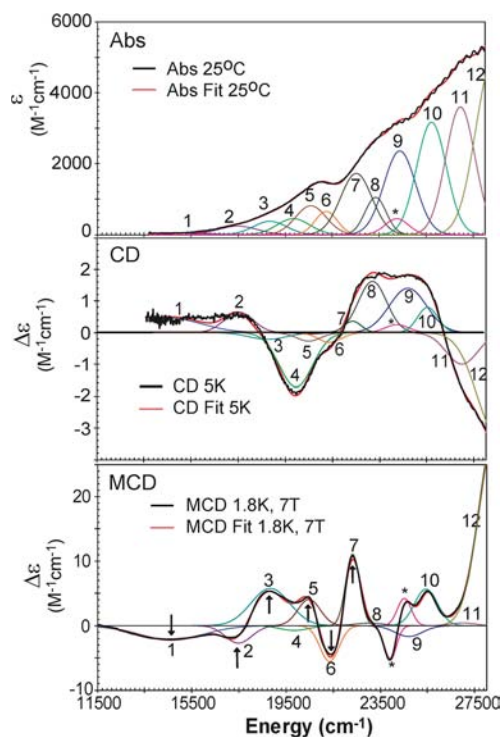


Figure 3. Abs, CD, and MCD spectra of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor in Ct RNR- β (R2). The arrows indicate the energies where the VTVH data were collected, and the * indicates a minor heme contaminant in the sample.

simultaneous Gaussian fit of the spectra in Figure 3 requires at least 12 bands in the range from 11500 cm^{-1} to 28000 cm^{-1} [note that there is a small contribution from a heme contaminant (*)]. These bands are assigned on the basis of their energies and their C_0/D_0 (MCD to Abs) intensity ratios. Because all d–d transitions are spin forbidden in a high-spin d^5 Fe^{III} site, any d–d transition with significant intensity in these spectra must involve the Mn^{IV} , as these transitions are spin-allowed. (A forbidden “spin-flip” transition for the d^5 Fe^{III} would be distinguished by its narrow bandwidth as there is no change in orbital occupancy). The C_0/D_0 ratio is given by

$$C_0/D_0 = k_B T / \mu_B B (\Delta\varepsilon/\varepsilon)_{\text{max}} \quad (1)$$

here k_B is the Boltzmann constant, μ_B is the Bohr magneton ($k_B/\mu_B = 1.489\text{ K}^{-1}\text{ T}$), T is the temperature in K, ε is the molar absorptivity in $\text{M}^{-1}\text{ cm}^{-1}$, and $\Delta\varepsilon$ is the MCD intensity maximum in $\text{M}^{-1}\text{ cm}^{-1}\text{ T}^{-1}$ recorded in the linear $1/T$ region. Features in the MCD spectra (Figure 3) exhibit approximately the same total intensity in left (positive) and right (negative) circular polarization, and therefore the data obey the sum rule. In this case, MCD intensity derives from spin–orbit coupling (SOC) between excited states (*vide infra*). Because the SOC constant for Mn^{IV} is much larger (415 cm^{-1}) than that of oxygen ($60\text{--}70\text{ cm}^{-1}$) or nitrogen (50 cm^{-1}) ligands,^{46,47} bands with high C_0/D_0 ratio have dominantly Mn character, which is the case for metal-centered LF transitions. Therefore, bands with high C_0/D_0 ratios must be Mn^{IV} d–d transitions, while bands with low C_0/D_0 can be assigned as ligand-to-metal CT (LMCT) transitions. From the C_0/D_0 ratios in Table 1, bands 1–4 are assigned as Mn^{IV} d–d transitions, while bands 5–12 are assigned as LMCT transitions. Because Mn^{IV} has a larger Z_{eff} than Fe^{III} ,^{48–50} CT transitions to Mn^{IV} should

Table 1. Simultaneous Abs/CD/MCD Peak Fits

band	Abs		CD			MCD			C ₀ /D ₀ ratio	assignment
	energy (cm ⁻¹)	FWHM	energy (cm ⁻¹)	FWHM	sign	energy (cm ⁻¹)	FWHM	sign		
1	14662	3663	14600	3205	+	14617	3206	-	0.0477	d-d
2	17480	1967	17520	1460	+	17500	1319	-	0.0146	d-d
3	18900	2003	18750	1883	-	18750	2002	+	0.0183	d-d
4	19850	1650	19951	1649	-	19950	1648	-	-	d-d
5	20594	1530	20480	1060	-	20500	1060	+	0.0074	CT
6	21307	1177	21380	1166	-	21380	1168	-	0.0038	CT
7	22540	1401	22355	871	+	22300	824	+	0.0091	CT
8	23400	1377	23195	1413	+	23200	1402	+	0.0002	CT
9	24490	1796	24700	1342	+	24700	1177	-	0.0018	CT
10	25725	1483	25535	1012	+	25500	1043	+	0.0028	CT
11	26959	1471	26917	1451	-	27000	1440	+	0.0001	CT
12	28354	1732	28604	2066	-	28393	1610	+	0.0080	CT

contribute at lower energy than CT transitions to Fe^{III}. However, a previous study on oxo-bridged Fe^{III} dimers showed that a bent μ -oxo can give rise to an intense oxo-to-Fe^{III} CT transition feature at ~ 19000 cm⁻¹,⁵¹ similar to the energy of the lowest-energy CT band (band 5) in Figure 3. Spectral assignments for the CT region are presented in section 3.2.4.

3.2.2. VTVH MCD. The ground-state spin-Hamiltonian parameters of the active site and the polarizations of each transition were obtained from VTVH MCD data. VTVH data were collected for different bands at the energies indicated by arrows in Figure 3 (MCD). VTVH data above 23000 cm⁻¹ could not be collected due to the decomposition of the cofactor induced by the high-energy photons (even at the low intensity exiting the monochromator). MCD saturation data were collected for the six bands at a series of temperatures with a range of fields and are plotted as relative intensity vs $\beta H/k_B T$ in Figure 4. The VTVH data were collected near the Gaussian maxima but at energies selected to minimize contributions from overlapping bands. For band 2, the contribution from band 3 was subtracted, as it has significant overlap. Bands 3 and 5 showed the same behavior, while all other bands exhibited different behaviors.

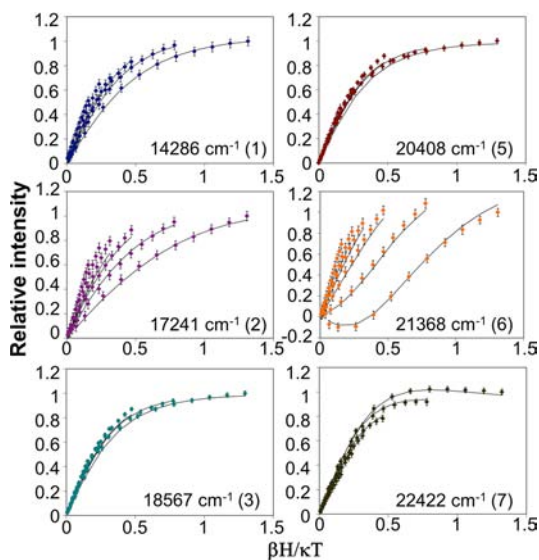


Figure 4. VTVH data for different bands, as indicated in Figure 3. Fits producing $|D| = 2.3$ cm⁻¹, $E/D = 0.33$ are shown as black lines.

These VTVH MCD data were fit using a spin-projection approach. This model correlates MCD intensity with the polarizations of transitions on the basis of the spin-expectation values of each of the sublevels in the dimer spin system. The relevant relation is⁵²

$$\frac{\Delta\epsilon}{E} = \frac{\gamma}{4\pi S} \int_0^\pi \int_0^{2\pi} \sum_i N_i \langle l_x \langle S_x \rangle_i M_{yz}^{\text{eff}} + l_y \langle S_y \rangle_i M_{xz}^{\text{eff}} + l_z \langle S_z \rangle_i M_{xy}^{\text{eff}} \rangle \sin\theta d\theta d\phi \quad (2)$$

In eq 2, x , y , and z refer to the principal axes of the ground-state spin Hamiltonian (*vide infra*); θ and ϕ are the polar angles between the incident light and the molecular z -axis; l_x , l_y , and l_z are the direction cosines for the magnetic field relative to the molecular coordinate system; M_{jk}^{eff} are the products of the polarizations of the electronic transitions; and γ is a collection of constants. This equation is also dependent on N_i , the Boltzmann population, and $\langle S_p \rangle_i$, the spin expectation value in the p direction for the spin sublevel i of the ground state. Two perpendicular transition moments, j and k , are required for MCD intensity. In a low-symmetry protein site, these transitions are unidirectional, so a nonzero M_{jk}^{eff} value is achieved by spin-orbit mixing with another excited state that has a perpendicularly polarized transition moment.

The Mn^{IV}Fe^{III} cofactor has a ground state of $S_{\text{tot}} = 1$ derived from antiferromagnetic coupling of Mn^{IV} ($S = 3/2$) with Fe^{III} ($S = 5/2$).²² An $S_{\text{tot}} = 1$ system has $M_s = 0$ and ± 1 sublevels, which are split by the axial zero-field splitting (ZFS, D). The $M_s = +1$ and -1 are split by rhombic splitting (E).

A series of D and E/D values were evaluated. The best-fit spin-Hamiltonian parameters are $|D| = \sim 2.3$ (± 0.4) cm⁻¹ and $E/D = 0.33$ (Figure 4). These values agree well with those determined independently from Mössbauer spectroscopy.²² The effective transition moment products M_{eff} were also obtained for each band. The percent polarization along each direction of the ZFS tensor for each transition can be determined from eq 3⁵² (% y , % z obtained by cyclic permutations of the indices.)

$$\%x = 100 \times \left\{ \frac{(M_{xy}^{\text{eff}} M_{xz}^{\text{eff}})^2}{(M_{xy}^{\text{eff}} M_{xz}^{\text{eff}})^2 + (M_{xy}^{\text{eff}} M_{yz}^{\text{eff}})^2 + (M_{yz}^{\text{eff}} M_{xz}^{\text{eff}})^2} \right\} \quad (3)$$

These polarization directions depend on the signs of D and E/D . However, regardless of the signs, the fits to the bands 5, 6, and 7 require three mutually perpendicular polarizations. In Section 3.2.4, the signs of D and E are both determined to be negative. The polarizations obtained for the different transitions with $D = -2.3 \text{ cm}^{-1}$ and $E/D = 0.33$ are given in Table 2.

Table 2. Effective Transition Moments and Percent Polarization for VTVH Data of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ Cofactor

band	M_{xy}	M_{xz}	M_{yz}	% x	% y	% z
1	0.44	1.13	0.50	40.2	7.9	51.9
2	0.45	2.30	-0.55	39.2	2.2	58.6
3 and 5	1.19	0.94	-0.25	89.7	6.3	4.0
6	-0.61	6.05	-2.23	6.9	0.9	92.2
7	1.14	-0.84	1.33	20.5	51.5	28.0

3.2.3. Ligand-Field (LF) Analysis. As presented in Section 2.1, the NRVS spectrum clearly requires the Mn^{IV} ion to be bound at site 1 (close to F_{127}). Allowing for different bridging (bis(μ -oxo) vs μ -oxo, μ -hydroxo) and protonation states of the site 1 solvent ligand (H_2O vs OH^-), four structural models for the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor was considered (note of these structures the one with a bis(μ -oxo) core and a terminal OH^- was also inconsistent with the NRVS data; see Section 3.1.2). The Mn^{IV} LF transition energies for these models were calculated by TD-DFT (see Figure S7c) and compared to the experimental LF transition energies for bands 1–4 assigned in Section 3.2.1. From the experimental energies in Figure 5e, the lowest-energy Mn^{IV} d–d transitions show a pattern in which bands 1 and 2 are split by $\sim 3000 \text{ cm}^{-1}$ and bands 2 and 3 are split by $\sim 1500 \text{ cm}^{-1}$. The calculated LF transitions for the four possible structural models from Section 2.1 are compared to the experimental splitting in Figure 5. The experimental pattern is consistent only with the μ -oxo, μ -hydroxo core structures for which these d–d transitions are calculated to have $3000\text{--}4000 \text{ cm}^{-1}$ (bands 1 and 2) and $1000\text{--}1500 \text{ cm}^{-1}$ (bands 2 and 3) splittings regardless of whether there is a terminal H_2O or OH^- ligand at the open coordination position at site 1. For structures with bis(μ -oxo) cores, the calculated LF transitions have the opposite splitting pattern, with transitions 1 and 2 split by $<1500 \text{ cm}^{-1}$, and bands 2 and 3 split by $\sim 3000 \text{ cm}^{-1}$. Therefore, bis(μ -oxo) core structures can be ruled out for the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor, consistent with the conclusion reached in a

previous EXAFS and DFT study that the Mn–Fe distance of 2.9 \AA is too long for a complex with two μ -oxo ligands.²⁷ This difference reflects the μ -oxo orbital mixing into the occupied Mn $d\pi$ orbitals. The structures with the oxo, hydroxo bridges have significant oxo- π out-of-plane orbital mixing into the Mn d_{yz} orbital from π overlap, only moderate oxo- π in-plane mixing into the d_{xz} orbital (because this in-plane orbital acquires some σ -character in the bent structure), and no oxo- π mixing into the d_{xy} orbital (z along the Mn- μ -oxo bond and y perpendicular to the bent Mn–O–Fe plane). These interactions split the energies of the occupied d orbitals and result in the large difference in the energies of transitions 1 ($d_{yz} \rightarrow d_{x^2-y^2}$) and 2 ($d_{xz} \rightarrow d_{x^2-y^2}$). In contrast, the structures with bis(μ -oxo) cores would have significant oxo- π out-of-plane mixing into both the d_{xz} and d_{yz} orbitals and less in-plane oxo- π mixing into $d_{x^2-y^2}$ (z perpendicular to the M_2O_2 plane, x along Mn–Fe). This pattern would increase the energies of both the d_{xz} and d_{yz} occupied orbitals relative to $d_{x^2-y^2}$ and result in a small calculated splitting between transitions 1 ($d_{yz} \rightarrow d_{z^2}$) and 2 ($d_{xz} \rightarrow d_{z^2}$) (Figure 5, right).

3.2.4. CT Analyses. **3.2.4.1. TD-DFT Calculations.** From the previous section, the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ active site is defined to have a μ -oxo, μ -hydroxo, (mono)- μ -carboxylato core structure. TD-DFT calculations were performed for the two remaining possible structures with this geometry (terminal OH^- vs H_2O on Mn^{IV}), and the LMCT transitions obtained were compared to the experimentally observed CT transitions. To obtain realistic estimates of the transition energies for comparison to the experimental values, CT to Mn^{IV} was calibrated according to a bis(μ -oxo) $\text{Mn}^{\text{IV}}\text{Mn}^{\text{III}}$ model complex,⁵³ and CT to Fe^{III} was calibrated according to bis(μ -oxo) and μ -oxo, μ -hydroxo $\text{Fe}^{\text{III}}\text{Fe}^{\text{III}}$ complexes.^{54–56} For each calibration, the TD-DFT calculations reproduced the CT transitions for the associated model complex but with some shift in the energies. (Figures S7 and S8: for CT to Mn^{IV} the energy shift is -5000 cm^{-1} ; for CT to Fe^{III} the shift is $+1500 \text{ cm}^{-1}$). In the calculations for the protein, for transitions localized on one metal site, the second metal site was replaced with Ga^{III} . Therefore, transitions to Mn^{IV} were calculated for a structure in which the Fe^{III} was substituted with Ga^{III} (geometry optimization resulted in only minor structural changes). Transitions to the Fe^{III} site were calculated for a structure in which the Mn^{IV} was changed to Ga^{III} (in this case, without further optimization to avoid structural changes due to the difference oxidation state; note

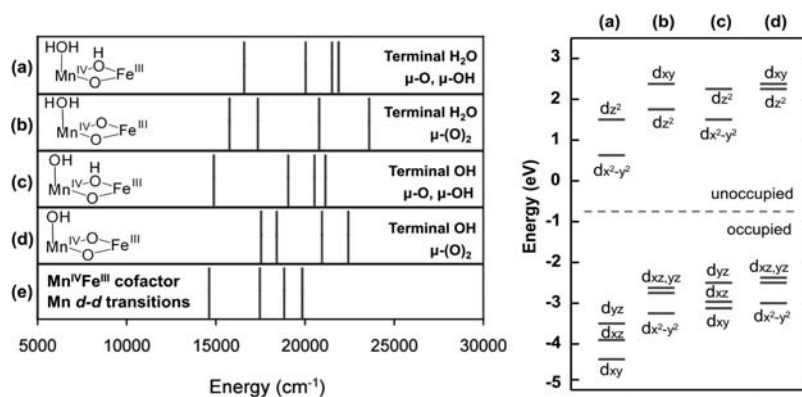


Figure 5. (Left) Comparison of the energies of the calculated d–d transitions (from TD-DFT) for four different structural models of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor (a–d) and the experimental d–d transitions from the MCD spectrum (e), and (right) relative Mn^{IV} d orbital energies for each structural model.

that further optimization with Ge^{IV} in place of Mn^{IV} still led to significant structural changes). The calculated LMCT transitions (energy shifted according to the calibrations) are shown in the bottom plots of Figure 6. For both structures (i.e., μ -oxo,

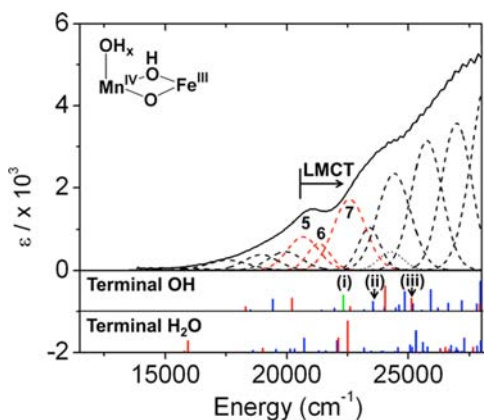


Figure 6. LMCT Abs spectra of Mn^{IV}Fe^{III} RNR- β (the three lowest CT bands are labeled bands 5–7) correlated to TD-DFT calculated μ -oxo \rightarrow Mn^{IV} (blue), terminal OH⁻ \rightarrow Mn^{IV} (green) and μ -oxo \rightarrow Fe^{III} (red) CT transitions for two possible structures, with either OH⁻ or H₂O as the terminal ligand on Mn at site 1. The orbitals for (i–iii) transitions are shown in Figure 7 (*vide infra*).

μ -hydroxo core with either a terminal OH⁻ or water on Mn at site 1), the calculated CT bands have high intensities starting at ~ 20000 cm⁻¹, consistent with the experimental Abs spectrum.⁵⁷ However, a noticeable difference is observed between the CT calculations for the two models. The model with terminal OH⁻ ligand has a low-energy CT transition to the Mn^{IV} (Figure 6, green bar), while the model with terminal H₂O does not. Therefore, the calculations show that a Mn^{IV}Fe^{III} complex with a terminal OH⁻ would have three types of LMCT transitions in the low-energy region: μ -oxo \rightarrow Mn^{IV}, terminal OH⁻ \rightarrow Mn^{IV}, and μ -oxo \rightarrow Fe^{III}, whereas a Mn^{IV}Fe^{III} structure with a terminal water would have only μ -oxo \rightarrow Mn^{IV} and Fe^{III} CT transitions. In the following section, this information is used to interpret CT transitions 5–7 and distinguish between these two possible models for the Mn^{IV}Fe^{III} cofactor.

3.2.4.2. Pseudo-A Terms. In the MCD spectrum, bands 5–7 show derivative shaped pseudo-A term behavior (Figure 3, bottom), and from VTVH MCD fitting, they were determined to have three perpendicular polarizations (Table 2).

For a spin-allowed transition from an orbitally non-degenerate $S_{\text{tot}} = 1$ ground state $|A\rangle$ to an excited state $|J\rangle$, the C_0 term intensity is given by⁵²

$$C_0(A \rightarrow J) = -\frac{1}{6} \sum_{uvw} g_w \sum_{K \neq A, J} [\Delta_{KJ}^{-1} L_w^{KJ} (D_u^{KA} D_v^{AJ} - D_v^{KA} D_u^{AJ}) + \Delta_{KA}^{-1} L_w^{KA} (D_u^{AJ} D_v^{JK} - D_v^{AJ} D_u^{JK})] \quad (4)$$

Here, g_w is the effective g -value of the $S = 1$ ground state in the w -direction, L_w^{KJ} and L_w^{KA} are the SOC matrix elements in the w -directions, and $D_u^{AJ} = \langle AS_{\text{tot}} M_s | m_u | JS_{\text{tot}} M_s \rangle (= D_u^{JA})$ is the component of the transition dipole moment between $|A\rangle$ and $|J\rangle$ in the u -direction (equivalent definitions for D_u^{KA} , D_v^{AJ} and D_v^{KA}).

For a unidirectional $|A\rangle \rightarrow |J\rangle$ transition, a nonzero C_0 value requires SOC between two excited states $|J\rangle$ and $|K\rangle$ of similar energies to which orthogonal polarized transitions are made

from the ground state $|A\rangle$. The resultant MCD features from this mechanism obey the sum rule (equal intensities in both circular polarizations) and show a derivative signal pattern (expressed by the first term of the eq 4). This mechanism involves SOC along the third mutually orthogonal direction, which can involve metal-centered SOC (from different ligands to one metal site) or ligand-centered SOC (from one ligand to two metal centers). In Table 2, there are three perpendicular CT-transition polarizations for bands 5–7, which also show pseudo-A term MCD behavior in Figure 3. This fact implies that there must be two sets of SOC among orthogonal transitions. These couplings can be between metal-centered oxo \rightarrow Mn^{IV} and terminal OH⁻ \rightarrow Mn^{IV} CT or between ligand-centered oxo \rightarrow Mn^{IV} and oxo \rightarrow Fe^{III} CT transitions. The Mn^{IV}Fe^{III} cluster will have three perpendicular polarizations only with a terminal OH⁻ on the Mn^{IV} site. Therefore, the Mn^{IV}Fe^{III} active site is defined as having a single bridging carboxylate, Mn^{IV} in site 1, μ -oxo, μ -hydroxo core, and a terminal OH on the Mn^{IV} ion.

Because Mn^{IV} has a larger Z_{eff} than Fe^{III}, the two lowest-energy transitions (band 5 and 6) should likely be assigned as terminal OH⁻ \rightarrow Mn^{IV} and bridging oxo \rightarrow Mn^{IV} CT. Band 7 would then be assigned as the bridging oxo \rightarrow Fe^{III} CT transition. From TD-DFT calculations, a possible assignment for bands 5–7 involving the pseudo-A intensity mechanism in eq 4 is given in Figure 7. Finally, these assignments are also suggested by the polarizations of the LF bands (bands 1–3, in Figure 3) obtained from the VTVH MCD analysis (Table 2). As d–d transitions borrow intensity from CT transitions and

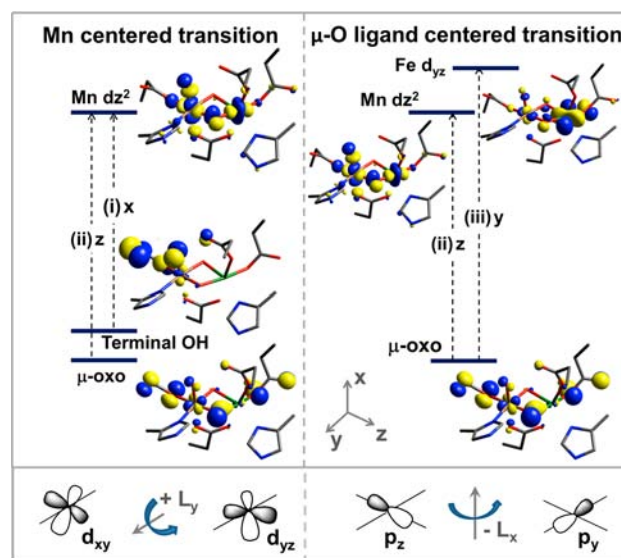


Figure 7. LMCT processes associated with pseudo-A term features. For pseudo-A term MCD behavior, there must be two perpendicular CT transitions and SOC along a third mutually orthogonal direction. SOC is a single center $1e^-$ operator. Therefore, two CTs must be “from” the same orbital or “to” the same orbital. Among the CT transitions in the TD-DFT calculations, there are three transitions (i–iii) involving orbitals which could SOC via a third perpendicular direction. For Mn-centered transitions, the terminal OH⁻ \rightarrow Mn CT transition is x -polarized, and the μ -oxo \rightarrow Mn CT transition is z -polarized. The ligand donor orbitals are mixed with d_{xy} and d_{yz} , which SOC via L_y (bottom left). For μ -oxo ligand-centered transitions, the CT to Mn is z -polarized and the transition to Fe is y -polarized. These metal d orbitals are mixed with the O_{pz} and O_{py} ligand orbitals, which SOC via L_x (bottom right).

are Mn^{IV}-centered, the polarizations for bands 1–3 should reflect the polarizations of the CT to the Mn^{IV}. From Table 2, bands 1–3 are *x*- and *z*-polarized, corresponding to the polarizations of transitions 5 (*x*-polarized) and 6 (*z*-polarized). They can thereby be assigned as terminal OH[−] and μ -oxo to Mn^{IV} CT transitions.

3.2.4.3. ZFS Tensor and Polarizations. CT transitions are polarized along the specific ligand–metal bond involved in the CT process, whereas the polarization directions determined by VTVH MCD are defined by the axes of the ZFS tensor. The coordinate frame of the ZFS tensor was calculated from the geometry-optimized structure with Mn^{IV} at site 1, a terminal carboxylate E₂₂₇ on Fe site 2, OH[−] on the Mn, and μ -oxo/ μ -OH core. This coordinate frame was then compared to the polarizations derived from VTVH MCD. The ZFS tensor for this binuclear site was calculated using⁵⁸

$$D_{\text{tot}(S=1)} = 3/10D_{\text{Mn(IV)}} + 14/5D_{\text{Fe(III)}} \quad (5)$$

by re-diagonalizing the combined ZFS matrix obtained for each metal center. The resultant tensor of the dimer is mostly governed by the ZFS of the Fe^{III} site. The strong contribution from the Fe^{III} is consistent with the rhombic behavior obtained from the fitting of the VTVH MCD and Mössbauer data,²² as the Fe^{III} center does not have a dominant, axial direction owing to the stronger bonding of the μ -oxo bridge to the Mn^{IV} ion. The orientations of the principal directions of the ZFS are shown in Figure 8 (blue dotted vectors). The bridging oxygen

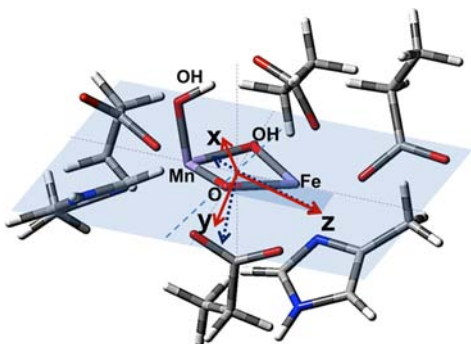


Figure 8. Calculated ZFS tensor orientation and intensity for the Mn^{IV}Fe^{III} cofactor. Initial ZFS orientation is shown with blue dotted arrows, where *x* and *y* orientations rotate to red arrows due to the antisymmetric ZFS arising from heterobimetallic center.

atoms are butterflyed $\sim 10^\circ$ out of the Mn-(μ -O, μ -OH)-Fe plane. The bond between the terminal OH[−] ligand and Mn is almost perpendicular to the Mn-(μ -O, μ -OH)-Fe plane. The *z*-axis of the ZFS tensor is in the Mn-(μ -O, μ -OH)-Fe plane, tilted 12° from the Mn-Fe vector, in the direction of the Mn-oxo bond. The *x*- and *y*-axes are rotated $+30^\circ$ and -60° , respectively, out of the Mn-(μ -O, μ -OH)-Fe plane (Figure 8). The Mn-ligand bonds have the following projections on this ZFS tensor: Mn-(μ -oxo): 81% *z*, 5% *x*, and 14% *y*; Mn-OH: 25% *x* and 75% *y*; Fe-(μ -oxo): 45% *z*, 41% *x*, and 14% *y*. The projection of the Mn-(μ -oxo) bond is consistent with the description of band 6, as *z*-polarized and its assignment as a μ -oxo \rightarrow Mn^{IV} CT transition. Conversely, the terminal OH[−] \rightarrow Mn^{IV} and μ -oxo \rightarrow Fe^{III} CT transitions have different projected polarizations relative to the experiment for bands 5 and 7. However, this structure should lead to an antisymmetric exchange contribution to the ZFS tensor, arising from the absence of inversion symmetry in the heterobimetallic center.

This antisymmetric exchange interaction should be on the order of $\sim 1 \text{ cm}^{-1}$ ($\sim (\Delta g/g)$)^{59,60} and in the direction perpendicular to the Mn-(μ -O, μ -OH)-Fe plane. This would rotate the *x* and *y* directions that have large out-of-plane components (Figure 8, solid vectors). Therefore, the projection of the OH[−] \rightarrow Mn^{IV} CT would increase in *x*-polarization, and the projection of the μ -oxo \rightarrow Fe^{III} CT would increase in *y*-polarization, leading to the assignments of their associated CT transitions as bands 5 and 7, respectively.

Table 3. Comparison of Calculated H-Atom Abstraction Thermodynamic Energies for Different Mn^{IV}Fe^{III} Structures

Reaction (H ⁺ , e [−])	Energy ⁽¹⁾ (kcal/mol)	
	Calc.	(Exp.) ⁶⁴
a.	-85.6	(-88.7 ± 0.6)
b.	-102.1	
c.	-83.1	
d.	-83.3	
e.	-82.5	
f.	-90.3	



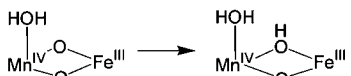
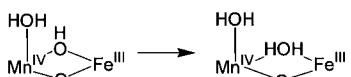
(1) energy corrected for H· atom energy with same functional and basis set (313.7 kcal/mol).

3.3. H-Atom Abstraction Energies and Proton Affinities. To evaluate the contributions of the Mn^{IV} position and the nature of the terminal ligand to the Mn^{IV}Fe^{III} cofactor site reactivity, the energies of H-atom abstraction reactions were calculated for the different possible structural models. Table 3 gives H-atom (H⁺, e[−]) abstraction energies (all energies used the B3LYP* functional for correlation to previous calculations)⁴⁴ for the class Ia Tyr and diiron centers, and the class Ic center with Mn at site 1 or 2 and either the OH[−] or H₂O at the open site 1.

The H-atom abstraction energy for the Fe^{III}Fe^{IV} complex, X, (Table 3, b) is larger by 12–20 kcal/mol than the H-atom abstraction energy for the Tyr and the four reactions of the Mn^{IV}Fe^{III} cofactor (a, c–f). The H-atom abstraction energy of Y• (a; −85.6 kcal/mol) was calculated to be similar to those of three of the Mn^{IV}Fe^{III} structures (c–e; −83.1, −83.3, and −82.5 kcal/mol, respectively).⁶⁴ However, when OH[−] is the terminal ligand at site 1, in contrast to the results of the earlier computational study, the position of the Mn^{IV} is found to have an impact, as the Fe^{III}Mn^{IV} structure (f) has a 7.8 kcal/mol greater H-atom abstraction energy than the Mn^{IV}Fe^{III} structure (e).

To evaluate the selection of the μ -oxo/ μ -hydroxo core structure over that with a bis(μ -oxo) core as the functioning model for the Ct RNR cofactor, the proton affinities of the four structures were evaluated with Mn^{IV} at site 1 and either an OH^- or H_2O ligand at this open site. The differences in the proton affinities of structures (1) and (2) and (3) and (4) in Table 4 are about 35 kcal/mol, mostly deriving from the charge

Table 4. Relative Proton Affinities among Structures with Mn^{IV} at Site 1

Reaction (H^+)	$\Delta\Delta E$ (kcal/mol)
(1) 	-35.9
(2) 	0
(3) 	-1.2
(4) 	33.9

difference between the pairs of sites. Thus the proton affinities for (2) and (3), where a proton is shifted from the bridge to the terminal hydroxide, are essentially the same.

These computational results are considered below in terms of the contribution of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ active site structure derived from NRVS/MCD spectroscopic studies to enzyme reactivity and its relevance to the class Ia radical-translocation mechanism.

4. DISCUSSION

We have considered 18 possible structures for the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor and used the combination of NRVS and Abs/CD/VTVH MCD spectroscopies and DFT calculations to define its geometric and electronic structure. The NRVS spectrum (Figure 1a) reflects the amount of Fe motion in the active site vibrations and exhibits a good correlation to the calculated spectra only when there is a single carboxylate bridge and Mn^{IV} is at site 1. This Mn^{IV} position is consistent with the results of recent X-ray crystallographic studies employing anomalous diffraction measurements.^{28–30} The terminal carboxylate, E_{227} , at Fe in site 2 provides an open coordination position at the Mn^{IV} in site 1 for solvent coordination. The Abs/CD/MCD/VTVH MCD spectra are dominated by the spin-allowed LF transitions, which are present only for Mn^{IV} , and the CT transitions that are able to produce MCD intensity through SOC. The LF energy-splitting pattern in Figure 5 (bands 1–4) requires a μ -oxo, μ -hydroxo core, and the pseudo-A term MCD behavior, and polarizations of the CT bands (band 5–7, Figures 3 and 4, Table 2) indicate that the terminal ligand at the exchangeable position on the Mn^{IV} is hydroxo rather than water. These results and analyses lead to the optimized $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ active site structure in Figure 9. It has 2 short Mn–O bonds from the μ -oxo and the terminal hydroxo ligand (~ 1.79 Å). This geometry correlates well with EXAFS data, which show 1 or 2 short Mn–O bonds at ~ 1.74 Å.²⁷ The

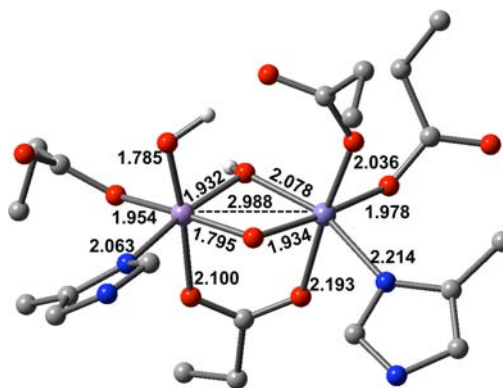


Figure 9. DFT-optimized structural model for the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor from this study.

calculated Mn–Fe distance is 2.99 Å, where EXAFS found a Mn–Fe distance of 2.92 Å. The calculated exchange coupling is $J = -97.3 \text{ cm}^{-1}$ ($-2J\text{S}_1\text{S}_2$ Hamiltonian) giving an $S_{\text{tot}} = 1$ ground state (from antiferromagnetic coupling of Mn^{IV} ($S = 3/2$) and Fe^{III} ($S = 5/2$)), consistent with Mössbauer spectroscopic results. The Mössbauer parameters calculated for this structure are $\delta = 0.63$ and $\Delta E_{\text{Q}} = -1.56 \text{ mm/s}$ ($\eta = 0.33$), which are in reasonable agreement with the experimental values of $\delta = 0.52 \text{ mm/s}$ and $\Delta E_{\text{Q}} = -1.32 \text{ mm/s}$ ($\eta = 0.11$).^{22,27} Therefore, the structure in Figure 9 is also consistent with existing spectroscopic data. These geometric and electronic structure of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ active site provide insight into its reactivity in radical translocation in the class Ic RNRs and its relation to the class Ia enzymes.

In class Ia RNRs, the $\text{Fe}^{\text{III}}\text{Fe}^{\text{IV}}$ intermediate X oxidizes the nearby tyrosine in the β (R2) subunit to generate the tyrosine radical that reversibly transfers its electron hole to the substrate site in the α (R1) subunit. However, in the class Ic Ct RNR, there is no nearby tyrosine in β (R2), and the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ active site itself plays the role of the Y• in electron hole transport. Therefore, the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor should have a free energy for uptake of a H-atom equivalent or similar to that of the tyrosine radical in *E. coli* RNR- β (R2). This matching of thermodynamics is best accomplished by having a Mn^{IV} rather than an Fe^{IV} in its active site (Table 3, structure e vs b), which depresses the H-atom abstraction energy to a value similar to that of the tyrosine radical (Table 3, a). This notion is consistent with the considerations in ref 44. However, unlike this earlier computational study,⁴⁴ which used a terminal water ligand at the Mn^{IV} site and found similar reaction energies for $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ and $\text{Fe}^{\text{III}}\text{Mn}^{\text{IV}}$ structures (consistent with Table 3), the spectroscopic results presented here show that OH^- is the terminal ligand at site 1. For this exogenous ligand at site 1, the $\text{Fe}^{\text{III}}\text{Mn}^{\text{IV}}$ active site has a 7.8 kcal/mol greater energy for uptake of a H-atom than the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ model. This depression of the energy for H• uptake in the favored structural model arises from the bonding of the strong terminal OH^- donor to the Mn^{IV} , which stabilizes the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ reactant state to a greater extent than the corresponding $\text{Fe}^{\text{III}}\text{Mn}^{\text{IV}}$ reactant state (with the terminal OH coordinated to the Fe^{III}), bringing the free energy for H• uptake into alignment with that of the tyrosine radical (Table 3). An $\text{Fe}^{\text{III}}\text{Mn}^{\text{IV}}$ active site would thus have reduced reversibility in transfer of the e^- hole to the α (R1) subunit and thus poor reactivity, as observed experimentally.²⁹ In addition, an $\text{Fe}^{\text{III}}\text{Mn}^{\text{IV}}$ site would be

more susceptible to irreversible reduction by an external electron donor, which would inactivate the enzyme.

Interestingly, the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ (and $\text{Fe}^{\text{III}}\text{Mn}^{\text{IV}}$) models with a terminal water at site 1 also have calculated H-atom abstraction energies similar to that of the Y• (Table 3, c and d vs a). However, the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ structure with a terminal water has ~34 kcal/mol lower proton affinity compared to that of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ structure with terminal OH^- (Table 4, (2) and (4)). If class Ic RNRs parallel the radical translocation reaction that occurs in class Ia RNRs, this difference in proton affinity would provide mechanistic insight into radical translocation in class Ic RNRs. In class Ia RNRs, the tyrosine radical transfers an electron hole to the α (R1) subunit, with the Y_{356} residue found to be involved in this process. Previous studies of class Ia RNRs suggested that a conformational change in the α (R1) subunit upon substrate binding induces proton transfer (PT) from the cofactor site to Y_{122} • to initiate a PCET process.^{4,5} It was determined that the proton is transferred from the water ligand bound to the Fe^{III} at site 1 to Y_{122} •.⁶¹ In class Ic RNRs, the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ site replaces Y• and directly transfers an electron hole to the α (R1) subunit, and the equivalent protein residue to Y_{356} in class Ia RNRs, Y_{338} , is also present and involved in the radical translocation process.^{62,63} Therefore, in class Ic RNR, α (R1) subunit and substrate binding could induce a conformational change leading to PT, in this case to the terminal OH^- on the Mn^{IV} at site 1 to initiate the radical translocation process. From a comparison of the structures, the terminal carboxylate ligand E_{89} that replaces proton-transfer-mediating carboxylate D_{84} in the class Ia RNR- β (R2) subunits, could be involved in this conformational change. Alternatively, a structure with a terminal H_2O at Mn^{IV} and a bis(μ -oxo) core, which would have a proton affinity similar to that of the structure in Figure 9 with a terminal OH^- and μ -oxo/ μ - OH^- core (Table 4, (2) and (3)), would likely lack a mechanism to deliver the proton to the most basic site, the oxo bridge. Also, a structure with bis(μ -oxo) core and a terminal OH^- (Table 4, (1)) might have too high a proton affinity to allow reversible allosteric regulation of proton transfer. However, it should be noted that at this point no spectroscopic evidence for the proton transfer in class Ic RNRs has been observed, and thus substrate binding in the α (R1) subunit could induce a different conformational change in class Ic RNRs. Further spectroscopic studies on the effects of the α (R1) subunit and substrate binding on this cofactor site are required to understand the radical translocation mechanism in class Ic RNRs.

In conclusion, the combination of NRVS and optical (Abs/CD/MCD/VTM MCD) spectroscopies with DFT calculations has afforded a detailed picture of the geometric and electronic structure of the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ active site in the class Ic Ct RNR. Comparison of the properties (reduction potentials and proton affinities) calculated for this spectroscopically elucidated structure and its possible structural isomers suggests that this structure is optimized for stability in the resting state and reversibility of its electron-hole generation in the α (R1) subunit to initiate catalysis.

■ ASSOCIATED CONTENT

Supporting Information

$\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ cofactor Mössbauer spectra used for NRVS and MCD studies, calculated NRVS spectra for all 18 structural models, field- and temperature-dependent Abs (25 °C, 5 K), CD (25 °C, 5 K), and MCD spectra for the $\text{Mn}^{\text{IV}}\text{Fe}^{\text{III}}$ and $\text{Fe}^{\text{III}}\text{Mn}^{\text{IV}}$ forms of Ct RNR- β (R2), calibration of the TD-DFT

calculations (Section 3.2.4.1), the structure used for the TD-DFT calculations, and complete ref 37. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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Notes

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